Entropy Rate Balance for Closed Systems

Example: An inventor claims that the device shown generates electricity at a rate of 100 kJ/s while *receiving* a heat transfer of energy at a rate of **250 kJ/s** at a temperature of **500 K**, *receiving* a second heat transfer at a rate of **350 kJ/ s** at **700 K**, and *discharging* energy by heat transfer at a rate of **500 kJ/s** at a temperature of **1000 K**. Each heat transfer is positive in the direction of the accompanying arrow. For operation at steady state, evaluate this claim.

Entropy Rate Balance for Closed Systems

▶ Applying an energy rate balance at steady state

$$
\frac{dE^0}{dt} = 0 = \dot{Q}_1 + \dot{Q}_2 - \dot{Q}_3 - \dot{W}_e
$$

Solving $\dot{W}_e = 250 \text{ kJ/s} + 350 \text{ kJ/s} - 500 \text{ kJ/s} = 100 \text{ kJ/s}$

The claim is in accord with the first law of thermodynamics.

▶ Applying an entropy rate balance at steady state

$$
\frac{dS}{dt} = 0 = \frac{\dot{Q}_1}{T_1} + \frac{\dot{Q}_2}{T_2} - \frac{\dot{Q}_3}{T_3} + \dot{\sigma}
$$

Solving
$$
\dot{\sigma} = -\left(\frac{250 \text{ kJ/s}}{500 \text{ K}} + \frac{350 \text{ kJ/s}}{700 \text{ K}} - \frac{500 \text{ kJ/s}}{1000 \text{ K}}\right)
$$

$$
\dot{\sigma} = -(0.5 + 0.5 - 0.5)\frac{\text{kJ/s}}{\text{K}} = -0.5\frac{\text{kJ/s}}{\text{K}}
$$

Since $\dot{\sigma}$ is negative, the claim is not in accord with the second law of thermodynamics and is therefore false.

- ► Like mass and energy, entropy can be transferred into or out of a control volume by streams of matter.
- ▶ Since this is the principal difference between the closed system and control volume entropy rate balances, the control volume form can be obtained by modifying the closed system form to account for such entropy transfer. The result is

$$
\frac{dS_{\text{cv}}}{dt} = \sum_{j} \frac{\dot{Q}_j}{T_j} + \sum_{i} \dot{m}_i s_i - \sum_{e} \dot{m}_e s_e + \dot{\sigma}_{\text{cv}}
$$
\nrate of rate of rate of entropy
\nentropy
\nchange transfer
\nproduction
\n
$$
\begin{array}{c}\n\frac{dS_{\text{cv}}}{dt} = \sum_{j} \dot{m}_j s_j - \sum_{e} \dot{m}_e s_e + \dot{\sigma}_{\text{cv}} \\
\text{rate of energy} = \text{entropy} \\
\text{entropy} = \text{entropy} \\
\text{transfer} = \text{production}\n\end{array}
$$
\n(Bq. 6.34)

where $\dot{m}_i s_i$ and $\dot{m}_e s_e$ account, respectively, for rates of entropy **transfer accompanying mass flow at inlets** *i* **and exits** *e***.**

►For control volumes at steady state, **Eq. 6.34** reduces to give

$$
0 = \sum_{j} \frac{\dot{Q}_j}{T_j} + \sum_{i} \dot{m}_i s_i - \sum_{e} \dot{m}_e s_e + \dot{\sigma}_{cv} \qquad (Eq. 6.36)
$$

► For a one-inlet, one-exit control volume at steady state, **Eq. 6.36** reduces to give

$$
0 = \sum_{j} \frac{\dot{Q}_j}{T_j} + \dot{m}(s_1 - s_2) + \dot{\sigma}_{cv}
$$
 (Eq. 6.37)

where 1 and 2 denote the inlet and exit, respectively, and *m***is the common mass flow rate at these locations.**

Example: Water vapor enters a valve at 0.7 bar, 280^oC and exits at 0.35 bar. (a) If the water vapor undergoes a throttling process, determine the rate of entropy production within the valve, in kJ/K per kg of water vapor flowing. **(b)** What is the source of entropy production in this case?

(a) For a throttling process, there is no significant heat transfer. Thus, **Eq. 6.37** reduces to

$$
0 = \sum_{j} \frac{\dot{Q}_{j}^{0}}{T_{j}} + \dot{m}(s_{1} - s_{2}) + \dot{\sigma}_{cv} \rightarrow 0 = \dot{m}(s_{1} - s_{2}) + \dot{\sigma}_{cv}
$$

Solving

From Table A-4, $h_1 = 3035.0$ kJ/kg, $s_1 = 8.3162$ kJ/kg·K.

For a throttling process, $h_2 = h_1$ (Eq. 4.22). Interpolating in Table A-4 at 0.35 bar and $h_2 = 3035.0 \text{ kJ/kg}$, $s_2 = 8.6295 \text{ kJ/kg·K}.$

Finally
$$
\frac{\dot{\sigma}_{cv}}{\dot{m}}
$$
 = (8.6295 – 8.3162) kJ/kg·K = **0.3133 kJ/kg·K**

(b) Selecting from the list of irreversibilities provided in **Sec. 5.3.1**, the source of the entropy production here is the *unrestrained* expansion to a lower pressure undergone by the water vapor.

Comment: The value of the entropy production for a single component such as the throttling valve considered here often does not have much significance by itself. The significance of the entropy production of any component is normally determined through comparison with the entropy production values of other components combined with that component to form an integrated system. Reducing irreversibilities of components with the highest entropy production rates may lead to improved thermodynamic performance of the integrated system.

Calculating Entropy Change

►The property data provided in **Tables A-2** through **A-18**, similar compilations for other substances, and numerous important relations among such properties are established using the *TdS* equations. When expressed on a unit mass basis, these equations are

$$
T ds = du + p dv
$$
 (Eq. 6.10a)

$$
T ds = dh - v dp
$$
 (Eq. 6.10b)

$$
T ds = du + p dv
$$
 (Eq. 6.10a)
\n $T ds = dh - v dp$ (Eq. 6.10b)

$$
T ds = du + p dv
$$
 (Eq. 6.10a)

$$
T ds = dh - v dp
$$
 (Eq. 6.10b)

$$
(\delta Q)_{int} = T dS
$$

$$
Q_{int} = \int_{1}^{2} T dS
$$
 (Eq. 6.23)

Calculating Entropy Change

- ▶ As an application, consider a change in phase from saturated liquid to saturated vapor at constant pressure.
- ►Since pressure is constant, **Eq. 6.10b** reduces to give

$$
ds = \frac{dh}{T}
$$

▶ Then, because temperature is also constant during the phase change $s_{\rm g} - s_{\rm f} = \frac{h_{\rm g} - h_{\rm f}}{T}$ **(Eq. 6.12)**

This relationship is applied in property tables for tabulating ($s_{\rm g} - s_{\rm f}$) from known values of ($h_{\rm g} - h_{\rm f}$).

Calculating Entropy Change

▶ For example, consider water vapor at 100^oC **(373.15 K).** From **Table A-2**, $(h_g - h_f) = 2257.1$ kJ/kg. **Thus**

(*s***g –** *s***^f) =** (**2257.1 kJ/kg)/373.15 K = 6.049 kJ/kg·K**

which agrees with the value from Table A-2, as expected.

►Next, the *TdS* **equations** are applied to two additional cases: substances modeled as incompressible and gases modeled as ideal gases.

Calculating Entropy Change of an Incompressible Substance

- ▶ The incompressible substance model assumes the specific volume is constant and specific internal energy depends solely on temperature: $u = u(T)$. Thus, $du = c(T)dT$, where *c* denotes specific heat.
- ►With these relations, **Eq. 6.10a** reduces to give

$$
ds = \frac{c(T)dT}{T} + \frac{p d\psi^0}{T} = \frac{c(T)dT}{T}
$$

▶ On integration, the change in specific entropy is

$$
s_2 - s_1 = \int_{T_1}^{T_2} \frac{c(T)}{T} dT
$$

▶ When the specific heat is constant

$$
s_2 - s_1 = c \ln \frac{T_2}{T_1}
$$
 (incompressible, constant c)

(Eq. 6.13)

▶ The ideal gas model assumes pressure, specific volume and temperature are related by $pv = RT$. Also, specific internal energy and specific enthalpy each depend solely on temperature: $u = u(T)$, $h = h(T)$, giving $du = c_n dT$ and $dh = c_p dT$, respectively.

►Using these relations and integrating, the *TdS* **equations** give, respectively

$$
ds = \frac{du}{T} + \frac{p}{T} dv
$$

\n
$$
ds = c_v(T)\frac{dT}{T} + R\frac{dv}{v}
$$

\n
$$
ds = c_v(T)\frac{dT}{T} + R\frac{dv}{v}
$$

\n
$$
ds = c_p(T)\frac{dT}{T} - R\frac{dp}{p}
$$

\n
$$
s(T_2, v_2) - s(T_1, v_1) = \int_{T_1}^{T_2} c_v(T)\frac{dT}{T} + R\ln\frac{v_2}{v_1}
$$

\n
$$
s(T_2, p_2) - s(T_1, p_1) = \int_{T_1}^{T_2} c_p(T)\frac{dT}{T} - R\ln\frac{p_2}{p_1}
$$

\n**(Eq. 6.17)**
\n
$$
Fq. 6.18)
$$

- ► Since these particular equations give entropy change on a unit of mass basis, the constant *R* is determined from $R = R/M$.
- \blacktriangleright Since c_v and c_p are functions of temperature for ideal gases, such functional relations are required to perform the integration of the first term on the right of **Eqs. 6.17** and **6.18**.
- ▶ For several gases modeled as ideal gases, including air, $CO₂$, CO, O₂, N₂, and water vapor, the evaluation of entropy change can be reduced to a convenient tabular approach using the variable *s***^o** defined by

$$
s^{\circ}(T) = \int_{T'}^{T} \frac{c_{p}(T)}{T} dT \qquad (\mathbf{Eq. 6.19})
$$

where *T*' **is an arbitrary reference temperature**.

►Using *s***^o**, the integral term of **Eq. 6.18** can be expressed as

$$
\int_{T_1}^{T_2} c_p \frac{dT}{T} = \int_{T'}^{T_2} c_p \frac{dT}{T} - \int_{T'}^{T_1} c_p \frac{dT}{T}
$$

= $s^{\circ}(T_2) - s^{\circ}(T_1)$

►Accordingly, **Eq. 6.18** becomes

$$
s(T_2, p_2) - s(T_1, p_1) = s^{\circ}(T_2) - s^{\circ}(T_1) - R \ln \frac{p_2}{p_1}
$$
 (Eq. 6.20a)

or on a per mole basis as

$$
\bar{s}(T_2, p_2) - \bar{s}(T_1, p_1) = \bar{s}^{\circ}(T_2) - \bar{s}^{\circ}(T_1) - \bar{R} \ln \frac{p_2}{p_1}
$$
 (Eq. 6.20b)

► For air, Tables A-22 and A-22E provide s° in units of kJ/ kg K and Btu/lb[.] R, respectively. For the other gases mentioned, **Tables A-23** and \overline{A} -23E provide \overline{s}° in units of kJ/kmol·K and Btu/lbmol·^oR, respectively.

Example: Determine the change in specific entropy, in kJ/ kg \cdot K, of air as an ideal gas undergoing a process from T_1 = **300 K**, $p_1 = 1$ bar to $T_2 = 1420$ K, $p_2 = 5$ bar.

▶ From Table A-22, we get $s^0_1 = 1.70203$ and $s^0_2 = 3.37901$, each in kJ/kg·K. Substituting into **Eq. 6.20a**

$$
s_2 - s_1 = (3.37901 - 1.70203) \frac{kJ}{kg \cdot K} - \left(\frac{8.314}{28.97}\right) \left(\frac{kJ}{kg \cdot K}\right) \ln\left(\frac{5 \text{ bar}}{1 \text{ bar}}\right) = 1.215 \frac{kJ}{kg \cdot K}
$$

Ideal Gas Properties of Air

►**Tables A-22** and **A-22E** provide

additional data for air modeled as an ideal gas. These values, denoted by p_r and v_r , refer only to two states having the same specific entropy. This case has important applications, and is shown in the figure.

 \blacktriangleright When $s_2 = s_1$, the following equation relates T_1 , T_2 , p_1 , and p_2

$$
\frac{p_2}{p_1} = \frac{p_r(T_2)}{p_r(T_1)} \quad (s_1 = s_2, \text{ air only}) \quad (Eq. 6.41)
$$

where $p_r(T)$ is read from Table A-22 or A-22E, as **appropriate**.

Ideal Gas Properties of Air

 \blacktriangleright When $s_2 = s_1$, the following equation relates T_1 , T_2 , v_1 , and $v₂$

$$
\frac{v_2}{v_1} = \frac{v_r(T_2)}{v_r(T_1)} \quad (s_1 = s_2, \text{ air only}) \quad (Eq. 6.42)
$$

where $v_{\rm r}(T)$ is read from Table A-22 or A-22E, as **appropriate**.

Ideal Gas Properties of Air

Entropy Change of an Ideal Gas Assuming Constant Specific Heats

 \blacktriangleright When the specific heats c_v and c_p are assumed constant, **Eqs. 6.17** and **6.18** reduce, respectively, to

$$
s(T_2, v_2) - s(T_1, v_1) = \int_{T_1}^{T_2} c_v(T) \frac{dT}{T} + R \ln \frac{v_2}{v_1}
$$
\n
$$
s(T_2, p_2) - s(T_1, p_1) = \int_{T_1}^{T_2} c_p(T) \frac{dT}{T} - R \ln \frac{p_2}{p_1}
$$
\n
$$
s(T_2, v_2) - s(T_1, v_1) = c_v \ln \frac{T_2}{T_1} + R \ln \frac{v_2}{v_1}
$$
\n
$$
s(T_2, p_2) - s(T_1, p_1) = c_p \ln \frac{T_2}{T_1} - R \ln \frac{p_2}{p_1}
$$

(Eq. 6.21) (Eq. 6.22)

▶ These expressions have many applications. In particular, they can be applied to develop relations among *T*, *p*, and *v* at two states having the same specific entropy as shown in the figure.

Entropy Change of an Ideal Gas Assuming Constant Specific Heats

 \triangleright Since $s_2 = s_1$, **Eqs. 6.21** and **6.22** become

$$
0 = c_p \ln \frac{T_2}{T_1} - R \ln \frac{p_2}{p_1}
$$

$$
0 = c_v \ln \frac{T_2}{T_1} + R \ln \frac{v_2}{v_1}
$$

where *k* **is the specific ratio** ▶ With the ideal gas relations

$$
c_p(T) = \frac{kR}{k-1} \qquad c_v(T) = \frac{R}{k-1}
$$

▶ These equations can be solved, respectively, to give

$$
\frac{T_2}{T_1} = \left(\frac{p_2}{p_1}\right)^{(k-1)/k} \qquad (s_1 = s_2, \text{ constant } k) \qquad (Eq. 6.43)
$$
\n
$$
\frac{T_2}{T_1} = \left(\frac{v_1}{v_2}\right)^{k-1} \qquad (s_1 = s_2, \text{ constant } k) \qquad (Eq. 6.44)
$$

 \blacktriangleright Eliminating the **temperature ratio gives**

$$
\frac{p_2}{p_1} = \left(\frac{v_1}{v_2}\right)^k \qquad (s_1 = s_2, \text{ constant } k) \qquad (Eq. 6.45)
$$

Example: Air undergoes a process from $T_1 = 620$ K, $p_1 = 12$ bar to a final state where $s_2 = s_1$, $p_2 = 1.4$ bar. Employing the ideal gas model, determine the final temperature T_2 , in K. Solve using **(a)** p_r data from Table A-22 and **(b)** a constant specific heat ratio *k* evaluated at 620 K from Table A-20: $k = 1.374$. Comment.

(a) With Eq. 6.41 and $p_r(T_1) = 18.36$ from Table A-22

$$
p_r(T_2) = p_r(T_1) \left(\frac{p_2}{p_1}\right) = 18.36 \left(\frac{1.4 \text{ bar}}{12 \text{ bar}}\right) = 2.142
$$

Interpolating in Table A-22, $T_2 = 339.7$ K

Ideal Gas Properties of Air

(b) With **Eq. 6.43**

$$
T_2 = T_1 \left(\frac{p_2}{p_1}\right)^{(k-1)/k} = 620 \text{ K} \left(\frac{1.4 \text{ bar}}{12 \text{ bar}}\right)^{(0.374/1.374)}
$$

 T_2 = 345.5 K

Comment: The approach of **(a)** accounts for variation of specific heat with temperature but the approach of **(b)** does not. With a *k* value more representative of the temperature interval, the value obtained in **(b)** using **Eq. 6.43** would be in better agreement with that obtained in **(a)** with **Eq. 6.41**.

Isentropic Turbine Efficiency

- ► **If the change in kinetic energy of flowing matter is negligible,** $\frac{1}{2}(V_1^2 - V_2^2)$ drops out.
- ▶ If the change in potential energy of flowing matter is negligible, $g(z_1 - z_2)$ drops out.
- \blacktriangleright If the heat transfer with surroundings is negligible, $\dot{\mathcal{Q}}_{\text{cv}}$ drops **out.**

$$
\frac{\dot{W}_{\text{cv}}}{\dot{m}} = h_1 - h_2
$$

where

the left side is work *developed* **per unit of mass flowing**.

Isentropic Turbine Efficiency

▶ For a turbine, the entropy rate balance reduces to

$$
0 = \sum_{j} \sum_{j} \frac{\partial f_j}{T_j} + \dot{m}(s_1 - s_2) + \dot{\sigma}_{cv}
$$

▶ If the heat transfer with surroundings is negligible, \dot{q} rops out.

$$
\frac{\dot{\sigma}_{cv}}{\dot{m}} = s_2 - s_1 \ge 0
$$